

1 -

a) $x = vt = \beta ct = \beta c \gamma Z$

knowing that $p = \gamma m \beta c \Leftrightarrow \frac{p}{m} = \gamma \beta c$

$\Rightarrow x = \frac{p}{m} Z (c)$ and from PDG $c \tau_{\pi^+} = 7,8 \text{ m}$

\hookrightarrow to put in S.I. units

$x = \frac{145 \text{ MeV}}{140 \text{ MeV}} 7,8 \text{ m} = 8,08 \text{ m}$

b) (possible answer)

\rightarrow Use magnetic field to deflect particles

\rightarrow Deflection depends of particle mass, charge, velocity and of B

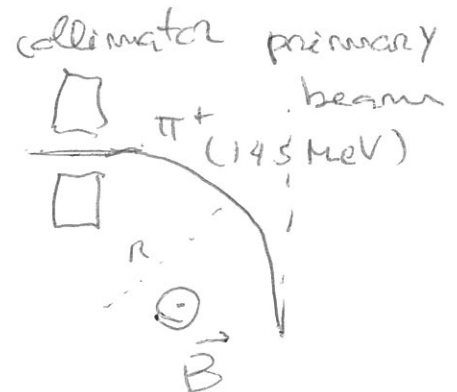
\rightarrow Use a collimator to stop close-by unwanted particle and create beam

\rightarrow Collimator has to be put before the π^+ decays ($x \gg \frac{2\pi r}{4}$)

$r_L = \frac{p_\pi}{qB}$; $r \sim \frac{2x}{\pi} \approx 5 \text{ m}$

$x \equiv$ mean ~~free~~ travelled distance

$\Rightarrow B = \frac{p_\pi}{2x |e|} \left(\frac{e}{c} \right) = 0,4 \text{ T}$
 \hookrightarrow S.I.



1-

$$c) \quad i) \quad \frac{dN}{dx} = - \underbrace{\frac{\sigma N_A \rho}{A}}_{\lambda} N \Rightarrow N = N_0 e^{-\lambda x}$$

$$\ln \left(\frac{N}{N_0} \right) = -\sigma \underbrace{\left(\frac{N_A \rho x}{A} \right)}_{y} = -\sigma y$$

projected number of protons $\equiv y$

$$\sigma = \frac{\ln(N_0/N)}{y} = 1,29 \times 10^{-25} \text{ cm}^2 = 129 \text{ mb}$$

where it was used

$$\frac{N_0}{N} = 1,116; \quad y = 8,5 \times 10^{23} \text{ cm}^{-2}; \quad 1 \text{ b} = 10^{-24} \text{ cm}^2$$

ii) From the PDG

$$\frac{dE}{dx} \sim 5 \text{ MeV} / (\text{g cm}^{-2})$$

$$E_{\text{loss}} = \frac{dE}{dx} [\text{MeV/g cm}^2] \cdot X [\text{g cm}^{-2}]$$

Knowing that the molar mass of hydrogen is

$$M_H = 1 \text{ g/mol} \quad \text{then,}$$

\$ 1) c) ii) (cont)

$$1 \text{ g} \longrightarrow 6,022 \times 10^{23} \text{ atoms} \equiv \text{mol}$$

$$m \longrightarrow 8,5 \times 10^{23} \text{ atoms}$$

$$\Rightarrow m = 1,41 \text{ g} \text{ and therefore } X = 1,41 \text{ g cm}^{-2} \text{ (depth)}$$

Hence,

$$E_{\text{loss}} = \frac{dE}{dX} \cdot X = 7,05 \text{ MeV}$$

~~From~~ for the dispersion from the PDG one has that the angle due to Coulomb multiple scattering is given ~~from~~ by,

$$\Theta_0 = \frac{13,6 \text{ MeV}}{BCP} z \sqrt{\frac{X}{X_0}} \left[1 + 0,038 \ln \left(\frac{X}{X_0} \right) \right]$$

$$X_0(\text{H}_2) = 63,04 \text{ g cm}^{-2} \text{ (end of PDG)}$$

$$X = 1,41 \text{ g cm}^{-2} ; C = 1 ; Z = 1$$

$$\Rightarrow \Theta_0 = 0,017 \text{ rad} \simeq 0,96^\circ$$

1-d) The angle that the μ^+ can have with the beam depends of the relation between β_μ^* and $\beta_{CM} \equiv \beta_\pi$

Two situations can be identified

$$\beta_\mu^* \geq \beta_{CM}$$

(LAB)



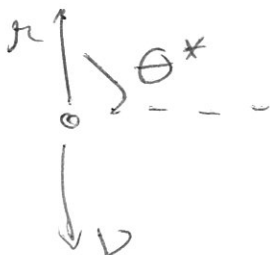
The μ is emitted backwards and the system boost is not big enough to reverse its velocity.

In this case the maximum possible angle with the beam is 180°

If the μ velocity (β_μ^*) in the CM is smaller than β_{CM} then the μ is always emitted in the forward direction.

In this case the μ maximum angle with the beam line occurs when the muon is emitted transversally to the boost (beam) direction. $\theta^* = \pi/2$

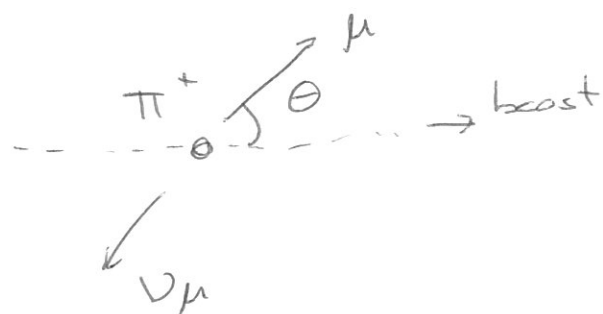
CM



LAB



1-d) In the CM



$$\gamma = \frac{E_\pi}{m_\pi} \quad \text{and} \quad E_\pi = \sqrt{p_\pi^2 + m_\pi^2}$$

$$\beta = \frac{p_\pi}{E_\pi}$$

First one needs to check if the muon can be emitted backwards in the LAB, i.e.,

β_μ^* is greater than $\beta_{CM} \equiv \beta_\pi$ or

$$\beta_\mu = \frac{p_\mu^*}{E_\mu^*} \quad \text{and} \quad \beta_{CM} = \frac{p_\pi}{\sqrt{p_\pi^2 + m_\pi^2}} = 0.72$$

To get p_μ^* (in the CM) one uses,

$$p_\pi = p_\nu + p_\mu \Leftrightarrow (4\text{-vectors})$$

$$(p_\pi - p_\mu)^2 = p_\nu^2 \Leftrightarrow p_\pi^2 + p_\mu^2 - 2p_\pi \cdot p_\mu = p_\nu^2 \quad (\text{eq. 1})$$

Knowing that in CM

$$p_\pi = (m_\pi, \vec{0})$$

$$p_\mu = (E_\mu^*, p_\mu^*)$$

$$p_\nu = (E_\nu^*, -p_\mu^*)$$

because

$$p_\mu = -p_\nu \quad \text{and}$$

$$E_\nu = p_\nu$$

Using the 4-vector inner product

$$p_\mu p_\mu = E^2 - p^2 \quad \text{and} \quad E^2 = p^2 + m^2$$

1-d) (cont.)

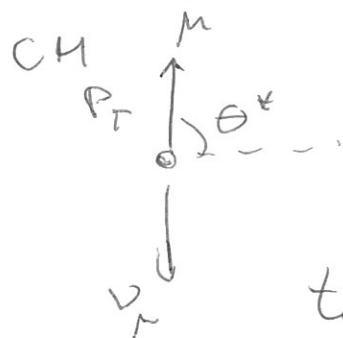
one can transform eq. 1 into

$$m_{\pi}^2 + m_{\mu}^2 - 2 E_{\mu}^* m_{\pi} = 0 \Leftrightarrow E_{\mu}^* = \frac{m_{\pi}^2 + m_{\mu}^2}{2 m_{\pi}}$$

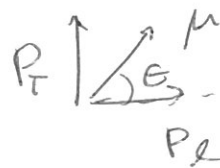
Hence, $E_{\mu}^* = 110 \text{ MeV}$ and

$$\beta_{\mu}^* = \frac{p_{\mu}^*}{E_{\mu}^*} = 0,298$$

Therefore $\beta_{\mu}^* < \beta_{CH}$ and the μ is always emitted in the direction of the boost. Thus, the maximum angle occurs for $\theta^* = 90^\circ$



LAB:



$$\tan \theta = \frac{p_T}{p_L}$$

$$p_{\mu}^* = p_T^* = p_T \quad (\text{no boost in the transverse direction})$$

$$\begin{pmatrix} E \\ p_L \end{pmatrix} = \begin{pmatrix} \gamma & \gamma \beta \\ \gamma \beta & \gamma \end{pmatrix} \begin{pmatrix} E^* \\ p_L^* \end{pmatrix}$$

$$\Rightarrow p_L = \gamma \beta E_{\mu}^* + \gamma p_{\mu}^* \cos \theta^* = \gamma \beta E_{\mu}^*$$

$$\tan \theta = \frac{p_{\mu}^*}{\gamma \beta E_{\mu}^*} = \frac{\sqrt{E_{\mu}^{*2} - m_{\mu}^2}}{\frac{E_{\pi}}{m_{\pi}} \frac{p_{\pi}}{E_{\pi}} E_{\mu}^*} \Rightarrow \theta \approx 16^\circ$$

1-d)

LAB energy

$$E_{\mu} = \gamma E_{\mu}^* + \gamma \beta \cos \theta^*$$

$$\text{Max } E: \theta^* = 0 \Rightarrow E_{\mu}^{\text{max}} = \gamma E_{\mu}^* + \gamma \beta P_{\mu}^* = 190.3 \text{ MeV}$$

$$\text{Min } E: \theta^* = 180^\circ \Rightarrow E_{\mu}^{\text{min}} = \gamma E_{\mu}^* - \gamma \beta P_{\mu}^* = 126.5 \text{ MeV}$$

knowing that in CM the muon angular production distribution is isotropic

$$\Rightarrow \frac{dN}{d\cos\theta^*} = \text{cte}$$

$$\frac{dN}{dE} = \frac{dN}{d\cos\theta^*} \frac{d\cos\theta^*}{dE}$$

\hookrightarrow energy spectrum in the LAB

$$\text{Since, } \cos\theta^* = \frac{E_{\mu} - \gamma E_{\mu}^*}{\gamma \beta P_{\mu}^*}$$

$$\frac{dN}{dE} = \text{cte} \times \frac{1}{\gamma \beta P_{\mu}^*} = \text{cte} \Rightarrow \text{uniform,}$$

Hence



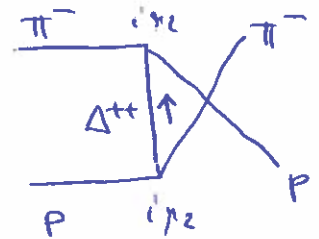
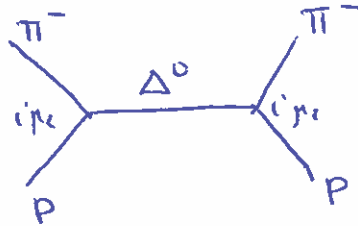
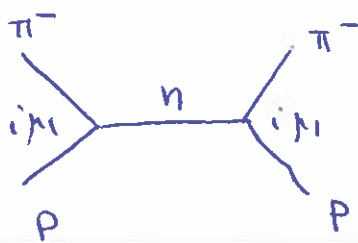
Problem 2

a) To be able to draw all diagrams we need to know the possible charges for the particles. The n is neutral and the p is positively charged. As for the π it is a triplet of isospin 1 with π^+, π^- and π^0 as you can look in the PDG.

The PDG also tells you that the Δ has $I = \frac{3}{2}$ and therefore $2 \times \frac{3}{2} + 1 = 4$ states. The PDG tells us that we have

$$\Delta^{++}, \Delta^+, \Delta^0, \Delta^-$$

With this knowledge we have the diagrams



where the indicated charges flow from left to right except the Δ^{++} that flows as indicated.

b) For the process to occur we should have in the CM frame

$$\sqrt{s} \geq m_p + m_{\Delta^0}$$

On the other hand the 4-momenta of the initial

(2)

particles in the Lab are

$$p_1 = (E_1, 0, 0, \sqrt{E_1^2 - m_n^2})$$

$$p_2 = (m_p, 0, 0, 0)$$

the Mandelstam variable $s = (p_1 + p_2)^2$ is an invariant and has the meaning of the square of the CM energy. therefore

$$s = (p_1 + p_2)^2 = m_p^2 + m_n^2 + 2E_1 m_p \geq (m_p + m_{\Delta^0})^2$$

or

$$2E_1 m_p = m_{\Delta^0}^2 + 2m_p m_{\Delta^0} - m_n^2$$

$$E_1 = \frac{m_{\Delta^0}^2 + 2m_p m_{\Delta^0} - m_n^2}{2m_p}$$

Numerically $m_p = 938.27 \text{ MeV}$

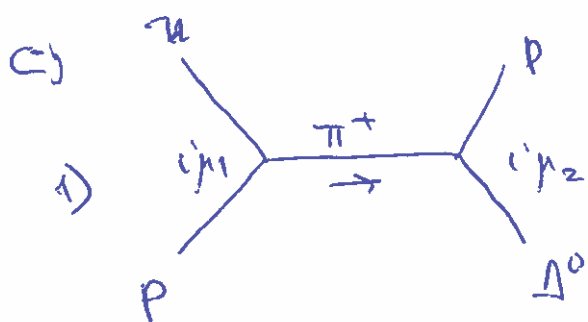
(PDG) $m_n = 939.56 \text{ MeV}$

$$m_{\Delta^0} = 1232 \text{ MeV}$$

and we get

$$E_1 = 1570.4 \text{ MeV}$$

(3)

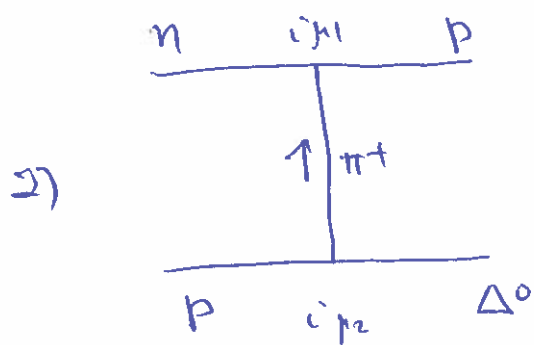


$$\Rightarrow -i M_1 = (i p_1)(i p_2) \frac{i}{(p_1 + p_2)^2 - m_\pi^2}$$

or

$$M_1 = \frac{p_1 p_2}{s - m_\pi^2}$$

$$s = (p_1 + p_2)^2$$

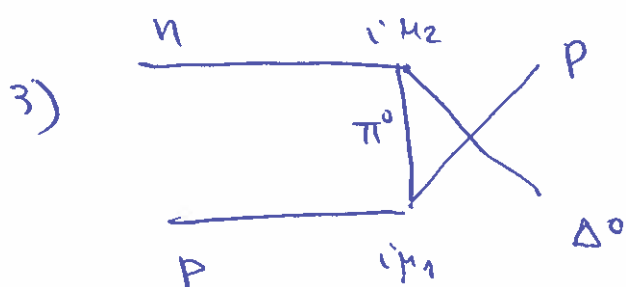


$$-i M_2 = (i p_1)(i p_2) \frac{i}{(p_1 - p_3)^2 - m_\pi^2}$$

or

$$M_2 = \frac{p_1 p_2}{t - m_\pi^2}$$

$$t = (p_1 - p_3)^2$$



$$-i M = (i p_1)(i p_2) \frac{i}{(p_1 - p_4)^2 - m_{\pi^0}^2}$$

or

$$M_3 = \frac{p_1 p_2}{u - m_{\pi^0}^2}$$

$$u = (p_1 - p_4)^2$$

(4)

a) We have

$$\frac{d\mathcal{T}}{d\Omega} = \frac{1}{64\pi^2 S} \frac{|\vec{p}_3^{CH}|}{|p_1^{CH}|} |M|^2$$

we have

$$E_1^{CH} = \frac{S + m_n^2 - m_p^2}{2\sqrt{S}}$$

$$E_3^{CH} = \frac{S + m_p^2 - m_{\Delta^0}^2}{2\sqrt{S}}$$

Therefore

$$|\vec{p}_1^{CH}| = \sqrt{(E_1^{CH})^2 - m_n^2} = \frac{1}{2\sqrt{S}} \sqrt{S^2 + m_n^4 + m_p^4 - 2Sm_n^2 - 2Sm_p^2 - 2m_n^2 m_p^2}$$

$$|\vec{p}_3^{CH}| = \sqrt{(E_3^{CH})^2 - m_p^2} = \frac{1}{2\sqrt{S}} \sqrt{S^2 + m_{\Delta^0}^4 + m_p^4 - 2Sm_p^2 - 2Sm_{\Delta^0}^2 - 2m_p^2 m_{\Delta^0}^2}$$

and

$$M = g_1 g_2 \left[\frac{1}{S - m_{\pi^+}^2} + \frac{1}{t - m_{\pi^+}^2} + \frac{1}{u - m_{\pi^0}^2} \right]$$

we get finally

$$\boxed{\frac{d\mathcal{T}}{d\Omega} = \frac{(g_1 g_2)^2}{64\pi^2 S} \frac{\sqrt{S^2 + m_p^4 + m_{\Delta^0}^4 - 2Sm_p^2 - 2Sm_{\Delta^0}^2 - 2m_p^2 m_{\Delta^0}^2}}{\sqrt{S + m_n^4 + m_p^4 - 2Sm_p^2 - 2Sm_n^2 - 2m_n^2 m_p^2}} \left[\frac{1}{S - m_{\pi^+}^2} + \frac{1}{t - m_{\pi^+}^2} + \frac{1}{u - m_{\pi^0}^2} \right]}$$

e) Limit $\sqrt{s} \gg$ all mass. this means

(5)

$$\sqrt{s} \gg m_\pi, m_p, m_\Delta$$

but also

$$t \gg m_{\pi^+}^2 \quad ; \quad u \gg m_{H^0}^2$$

we get

$$\frac{d\sigma}{d\Omega} = \frac{(\mu_1 \mu_2)^2}{64\pi^2 s} \left[\frac{1}{s} + \frac{1}{t} + \frac{1}{u} \right]^2$$

Now, t and u have angle, neglecting masses

$$t = (p_1 - p_3)^2 = -2p_1 \cdot p_3 = -2\left(\frac{\sqrt{s}}{2}\right)^2 (1 - \cos\theta)$$

$$u = (p_1 - p_4)^2 = -2p_1 \cdot p_4 = -2\left(\frac{\sqrt{s}}{2}\right)^2 (1 + \cos\theta)$$

where we have used for $\sqrt{s} \gg$ mass

$$p_1^{CM} = \frac{\sqrt{s}}{2} (1, 0, 0, 1)$$

$$p_3^{CM} = \frac{\sqrt{s}}{2} (1, 0, 0, \cos\theta)$$

$$p_4^{CM} = \frac{\sqrt{s}}{2} (1, 0, 0, -\cos\theta)$$

Now the integrations on the angles will cancel (6) without the mass. This is not a problem in real experiments because θ can not be 0, or π (these angles are inside the pipe). But in order to simplify the problem was asked only to evaluate the contribution from S-channel there are no angles and the integration in $d\Omega$ from $-\pi$ to π . We set therefore for that case

$$\sigma = \frac{(M_1 p_2)^2}{64\pi^2 S} \frac{1}{S^2} \times 4\pi$$

$$\boxed{\sigma = \frac{(M_1 p_2)^2}{16\pi S^3}}$$

As $[S] = \text{Mass}^2$ and $[p_1] = [p_2] = \text{Mass}$ we have

$$[\sigma] = \frac{(\text{Mass})^4}{(\text{Mass}^2)^3} = \frac{1}{\text{Mass}^2} \quad \text{as required.}$$